

**UNIVERSITY COLLEGE LONDON**

*University of London*

**EXAMINATION FOR INTERNAL STUDENTS**

*For The Following Qualifications:-*

*B.Sc. M.Sci.*

**Physics 2B28: Statistical Thermodynamics and Condensed Matter Physics**

**COURSE CODE : PHYS2B28**

**UNIT VALUE : 0.50**

**DATE : 28-MAY-03**

**TIME : 10.00**

**TIME ALLOWED : 2 Hours 30 Minutes**

Answer ALL SIX questions from section A and THREE questions from section B.

The numbers in square brackets in the right-hand margin indicate the provisional allocation of marks per sub-section of a question.

Boltzmann's constant is  $k$

Internal Energy is  $E$  and  $\beta = 1/kT$

Planck's constant is  $h$  and  $\hbar = \frac{h}{2\pi}$

Avogadro constant  $N_A = 6.02 \times 10^{23} \text{mole}^{-1}$

Elementary charge  $e = 1.6 \times 10^{-19} \text{C}$

All other thermodynamic symbols have their usual meaning

The Stirling approximation is  $\ln n! = n \ln n - n$

Single particle density of momentum states  $f(p)dp = \left(\frac{V}{h^3}\right) 4\pi p^2 dp$

### SECTION A

[Part marks]

1. A system is in contact with a large heat reservoir such that the pressure is constant throughout, and the final temperature following a process is the same as the initial temperature. Show that the maximum **useful** work done by the system  $W_u$ , is given by  $W_u \leq -\Delta G$ , where  $G$ , the Gibbs free energy is given by  $G = E - TS + pV$ . [4]

In an electrolysis cell, the overall reaction is the dissociation of a water molecule with the transfer of two electrons across the cell. If the Gibbs free energy of the process is  $+237 \text{ kJ mol}^{-1}$  determine the minimum EMF that must be applied to the cell for electrolysis to start, showing details of your reasoning. [3]

2. A rubber band in thermal equilibrium may be modelled by  $N$  freely pivoted inextensible links each of length  $d$ , which can point forwards or backwards along the band.

If a band consists of 14 links of which 10 point forwards along the band, identify the macrostate, and determine the number of accessible microstates (or the statistical weight). [3]

Use the fundamental equation of thermodynamics and the Boltzmann hypothesis to **qualitatively** explain the origin of the force required to extend the band. [3]

3. Write down the Boltzmann distribution of the probability  $p_r$  for a system to be in the  $r^{\text{th}}$  microstate with energy  $E_r$ . [2]

If the energy levels of a hydrogen atom are given by  $E_n = -\alpha/n^2 \text{ eV}$  and the degeneracy of the  $n^{\text{th}}$  level is  $2n^2$  show that the probability of hydrogen atoms in thermal equilibrium, to be in the  $n^{\text{th}}$  level, is proportional to

$$2n^2 e^{3\alpha/2n^2 \bar{E}}$$

where  $\bar{E}$  is the average kinetic energy of the atoms. [3]

What is the ratio of atoms in the state  $n = 3$  to those in the state  $n = 1$  for hydrogen atoms with mean energy  $\bar{E}$  of 1 eV in a stellar gas, if  $\alpha$  is 13.6 eV? [1]

4. Show that allowed values of the momentum components  $p_x$  and  $p_y$  of a particle confined to move on a square surface of side  $L$  are given by

$$p_x = \frac{h}{2L}n_x \text{ and } p_y = \frac{h}{2L}n_y,$$

where  $n_x$  and  $n_y$  are positive integers. [3]

Obtain the single particle density of momentum states in two dimensions. [3]

5. Explain briefly what is meant by the term **photon** in the context of black-body radiation and the term **phonon** in the context of a crystalline solid. Why can such systems correctly be treated by the Boltzmann distribution? [3]

Show that the partition function  $Z_r$  and the occupation number  $n_r$  of the  $r^{\text{th}}$  microstate for photons are given by

$$Z_r = \frac{1}{1 - e^{-\beta\epsilon_r}} \text{ and } n_r = -\frac{1}{\beta} \frac{\partial}{\partial \epsilon_r} (\ln Z_r) = \frac{1}{e^{\beta\epsilon_r} - 1},$$

where  $\epsilon_r = \hbar\omega_r$  is the energy and  $\omega_r$  is the angular frequency of the  $r^{\text{th}}$  microstate. [2]

Write down an expression, without simplification, for the energy density  $u(\omega_r, T)$  of black body radiation, clearly identifying the component parts of your expression. [3]

6. A system containing a quantal gas, is in contact with a heat bath with which it can exchange particles as well as heat. Show that the Gibbs distribution (or the grand canonical distribution) for a state with  $N_r$  particles,  $N_r = \sum_i n_i$  where the  $n_i$  are the occupation numbers, and an energy  $E_{N_r}$ , where  $E_{N_r} = \sum_i n_i \epsilon_i$ , factorises into the product of probabilities

$$p_i(n_i) = \frac{\exp(\beta(\mu - \epsilon_i)n_i)}{Z_{Gi}}$$

where  $\mu$  is the chemical potential. [4]

Show that for quantal statistics, the mean occupation number is given by

$$\bar{n}_i = kT \frac{\partial \ln Z_{Gi}}{\partial \mu}$$

and obtain  $\bar{n}_i$  for Fermi-Dirac statistics. [3]

## SECTION B

7. Starting with the fundamental equation of thermodynamics and using the third Maxwell equation  $\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial p}{\partial T}\right)_V$ , show that

$$\left(\frac{\partial E}{\partial V}\right)_T = \left(T \left(\frac{\partial p}{\partial T}\right)_V - p\right). \quad [4]$$

A mole of gas obeys the van der Waals equation  $\left(p + \frac{a}{V^2}\right)(V - b) = RT$  where  $a$  and  $b$  are constants. Show that  $\left(\frac{\partial E}{\partial V}\right)_T = \frac{a}{V^2}$ . [3]

By considering  $E(T, V)$  obtain an expression for  $E$ , the internal energy of a mole of van der Waals gas. [4]

By considering  $S(T, V)$  show that

$$dS = \frac{C_V}{T}dT + \frac{R}{V-b}dV$$

and obtain an expression for the entropy. [4]

Obtain the relation between the temperature and volume for the ideal adiabatic expansion of the van der Waals gas. [5]

8. Write down the Boltzmann equation relating the entropy  $S$  to the number of microstates (or statistical weight)  $\Omega$  of an isolated system, and state the condition of equilibrium. [3]

An isolated system of volume  $V$ , containing  $N$  particles is partitioned by a diathermal wall into two subsystems 1 and 2. Derive the equilibrium condition

$$\left(\frac{\partial S_1}{\partial E_1}\right)_{N,V} = \left(\frac{\partial S_2}{\partial E_2}\right)_{N,V}, \text{ and justify the relation } \frac{1}{T} = \left(\frac{\partial S}{\partial E}\right)_{N,V}. \quad [4]$$

Spin 1/2 magnetic dipoles in a paramagnetic solid align in a magnetic field  $B$  such that dipoles aligned parallel to the field have energy  $-\mu B$  and those antiparallel have energy  $+\mu B$ , where  $\mu$  is the dipole moment. Determine the number of microstates of the macrostate of  $N$  dipoles, with  $n$  dipoles aligned parallel to the field, and show using the above equilibrium condition that

$$n = \frac{N}{1 + e^{-x}} \text{ and } E = N\mu B \left(1 - \frac{2}{1 + e^{-x}}\right) \text{ where } x = \frac{2\mu B}{kT}.$$

Comment on the physical interpretation of the internal energy  $E$  at  $x = 0$  and  $x = \infty$ . [8]

Suppose now that the paramagnetic solid is made of spin  $J$  particles where the magnetic interaction energy  $-\mu \cdot \mathbf{B}$  is given by  $-g_J \mu_B m_J B$ , where  $g_J$  is the Lande  $g$  factor,  $\mu_B$  the Bohr magneton and  $m_J, -J \leq m_J \leq J$ , the magnetic quantum number. Explain why this system is treated by the Boltzmann distribution and obtain the partition function. [5]

9. By considering a system containing  $N$  particles and of volume  $V$  in contact through diathermal walls with a large heat bath, show that the probability  $p_r$  of the system being in its  $r^{\text{th}}$  microstate with energy  $E_r$  is given by the Boltzmann distribution [4]

$$p_r = \frac{e^{-\beta E_r}}{Z}, \text{ where the partition function } Z = \sum_r e^{-\beta E_r}.$$

The DNA molecule may be modelled as a zip fastener where one end of the fastener is always closed. From the open end,  $s$  links are open with total energy  $s\epsilon$  where  $\epsilon$  is the energy of an open link, and  $N - s$  links are closed. Show that the partition function  $Z$  is given by

$$Z = \frac{1 - x^N}{1 - x} \text{ where } x = e^{-\beta\epsilon}. \quad [5]$$

Prove that the mean number of open links  $\bar{s}$  is given by

$$\bar{s} = x \frac{\partial \ln Z}{\partial x} = N \frac{x^N}{x^N - 1} - \frac{x}{x - 1} \quad [6]$$

By considering the behaviour of  $\ln Z$  as  $x \rightarrow 1$  show the  $\bar{s} = N/2$  when  $x = 1$ . Under what conditions of degeneracy can  $x = 1$  be reached? [5]

10. For Bose-Einstein statistics the mean occupation number  $\bar{n}_r$  of a single particle state of energy  $\epsilon_r$  is given by

$$\bar{n}_r = \frac{1}{(e^{\beta(\epsilon_r - \mu)} - 1)}, \text{ and for the lowest energy state } \bar{n}_1 = \frac{1}{(Be^{\beta\epsilon_1} - 1)}$$

where  $\epsilon_1$  is extremely close to zero energy and  $B = e^{-\beta\mu}$ . Explain why  $B$  is expected to be just greater than one at low temperatures. [4]

Using the density of states factor in terms of the momentum, and the integral given at the end of this question, show that at high temperature  $B$  is proportional to  $T^{\frac{3}{2}}$  for  $N$  bosons in a volume  $V$ . [5]

Show that the appropriate density of states factor for  $N$  non-relativistic spinless ideal bosons in a volume  $V$ , each of mass  $m$  and kinetic energy  $\epsilon$  is given by

$$f(\epsilon)d\epsilon = 2\pi \frac{V}{h^3} (2m)^{\frac{3}{2}} \epsilon^{\frac{1}{2}} d\epsilon. \quad [5]$$

Obtain an expression for the internal energy  $E$  of an ideal boson system below  $T_c$  where  $T_c$  is the temperature at which  $B$  becomes just greater than one and  $\mu$  close to zero. [4]

What is the temperature dependence of the heat capacity  $C_V$  of the above system where  $T < T_c$ ? [2]

In this question you may use

$$\int_0^\infty x^2 e^{-\alpha x^2} dx = \frac{1}{4} \left( \frac{\pi}{\alpha^3} \right)^{\frac{1}{2}}.$$

11. The energy levels  $E_n$  of a particle of mass  $m$  moving in a one dimensional box of length  $L$  are given by

$$E_n = \frac{\hbar^2 \pi^2 n^2}{2mL^2}$$

Show that the partition function of a particle in a three dimensional cubic box of side  $L$ , assuming that summations are replaced by integrals and that  $\int_0^\infty e^{-\alpha x^2} dx = \frac{1}{2} \left( \frac{\pi}{\alpha} \right)^{\frac{1}{2}}$ , is given by

$$Z(1, T) = \left( \frac{mkT}{2\pi\hbar^2} \right)^{3/2} L^3 \quad [5]$$

Write down the classical partition function for a monatomic gas of  $N$  identical particles, explaining the factors in the expression. [3]

Assuming, without proof that the above yields the Sackur-Tetrode equation for the entropy  $S$ ,

$$S = Nk \left[ \ln \frac{V}{N} + \frac{3}{2} \ln \left( \frac{mkT}{2\pi\hbar^2} \right) + \frac{5}{2} \right]$$

use  $\mu = -T \left( \frac{\partial S}{\partial N} \right)_V$  to show that the chemical potential  $\mu$  is given by

$$\mu = kT \ln \frac{n}{n_Q} \quad \text{where } n = \frac{N}{V} \text{ and } n_Q = \left( \frac{mkT}{2\pi\hbar^2} \right)^{\frac{3}{2}}. \quad [5]$$

Oxygen circulating in the blood stream is absorbed by the myoglobin ( $Mb$ ) molecule to form  $MbO_2$  which has an energy  $\epsilon$  with respect to the  $Mb$  molecule. Assuming that the oxygen in the blood can be treated as an ideal classical gas, and that the grand canonical partition function is of the usual form  $Z_{Gi} = \sum_{n_i} e^{\beta(\mu - \epsilon_i)n_i}$ , show that the ratio  $f$  of the oxygenated molecules to those unoxxygenated molecules is given by

$$f = \frac{p}{(n_Q kT e^{\epsilon/kT} + p)} \quad [7]$$

where  $p$  is the partial pressure of the oxygen.